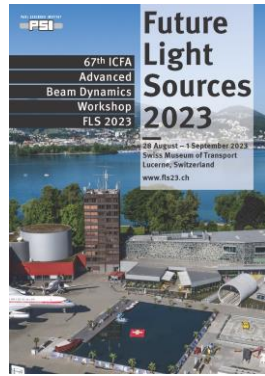




Elettra  
Sincrotrone  
Trieste



# Scaling of Beam Collective Effects with Bunch Charge in the CompactLight FEL

Simone Di Mitri

*Elettra Sincrotrone Trieste & University of Trieste, Dept. Physics*

on behalf of the CompactLight Collaboration



## Intro & Motivations

Space Charge-dominated Emittance

Coherent Synchrotron Radiation

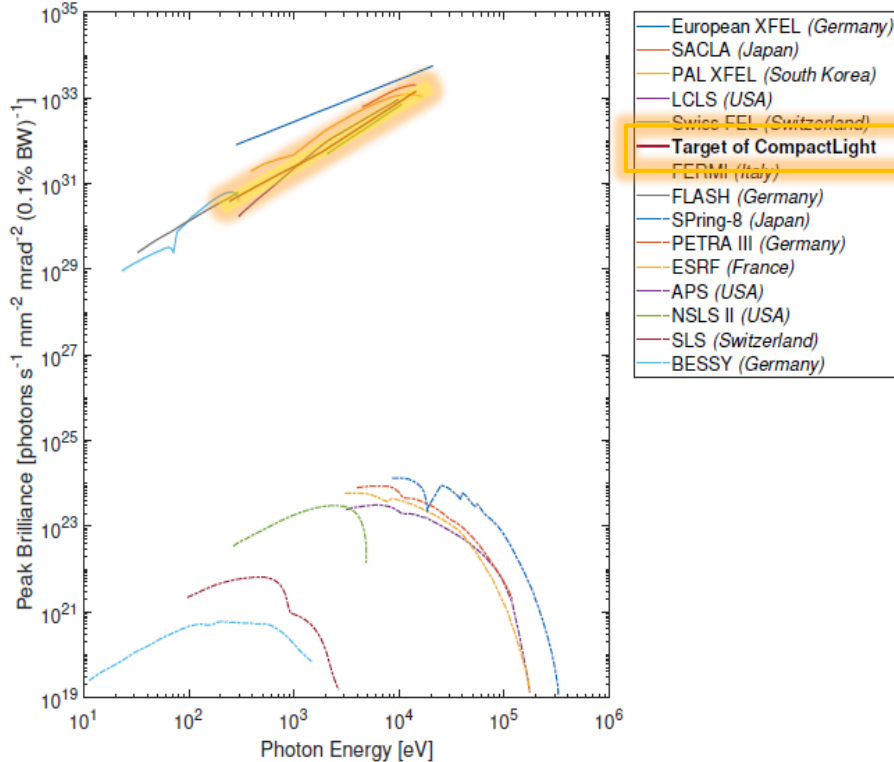
Beam Break-Up

FEL energy and transverse coherence

## Conclusions



# Introduction



$$\left\{ \begin{array}{l} \lambda_R \sim 1 \text{ \AA} \text{ (16 keV)} \\ \lambda_U = 10 \text{ mm, } K \sim 1 \text{ (SCU+AB)} \end{array} \right.$$

$$\hookrightarrow E_e = m_e c^2 \sqrt{\frac{\lambda_U}{2\lambda_R} \left(1 + \frac{K^2}{2}\right)} = 5.5 \text{ GeV}$$

1 kHz NC, 65 MV/m

$$\hookrightarrow L_{linac} \approx \frac{E_e}{G} = 100 \text{ m active length}$$

$$I \approx \frac{75 \text{ pC}}{\sqrt{2\pi} 30 \text{ fs}} = 1 \text{ kA}, \quad \rho \approx 0.01\%$$

$$\hookrightarrow E_{sase} \approx 2.4 \times \sigma_{t,b} \times \rho I E < 100 \mu\text{J}$$



- ❑ Scientific cases for several 100's  $\mu\text{J}$  in soft X-rays
  - Keep **same final peak current** to keep  $L_{sat}$  and  $P_{sat}$  almost fixed
    - 0.4 kA @ 1 kHz (Soft-X), 4.5 kA @ 0.1 kHz (Hard-X)
  - Bunch duration is increased proportionally to the bunch charge
  - We expect  $E_{sase} \propto Q$
  
- ❑ *But:* collective effects are also  $\propto Q_b$ , enlarging the projected e-beam emittances
  - In reality  $E_{sase} \propto Q^\nu$ ,  $\nu < 1$
  
- ❑ **Goal: set up an analytical model to estimate  $E_{sase}$  and  $F_{coh}$  vs.  $Q$  ( $< 1$  nC)**



# Space Charge-dominated Emittance

□ Envelope eq., cylindrical beam,  $\varepsilon_{th} \rightarrow 0$ :

[1] J. Rosenzweig and E. Colby, TESLA note 95-04

$$\sigma_x'' + \sigma_x' \frac{(\beta\gamma)'}{\beta\gamma} + K\sigma_x = \kappa_{sc}\sigma_x$$

with  $\kappa_{sc} := \frac{2I}{I_A} \frac{1}{(\beta\gamma)^3 \sigma_x^2} f \left( \frac{\sigma_x}{\beta\gamma\sigma_z} \right) = \left( \frac{2c}{I_A} \frac{f/g}{\beta^2\gamma^3} \right) \frac{Q}{\sigma_z\sigma_x^2} \equiv \text{const.}$  for increasing Q

▪ Keep the aspect ratio **f** and the longitudinal profile **g** constant  $\Rightarrow$

$$\frac{Q}{\sigma_z\sigma_x^2} \equiv \text{const.}$$

□ For linear SC (blow-out),  $\sigma_i \propto Q^{1/3}$ . Since  $\sigma_x = \sqrt{\varepsilon_x\beta_x}$  and  $\beta_x \approx \text{const.}$ , we find:

$$\varepsilon_x [\mu\text{m}] \approx a \times Q [\text{nC}]^{2/3}$$

and

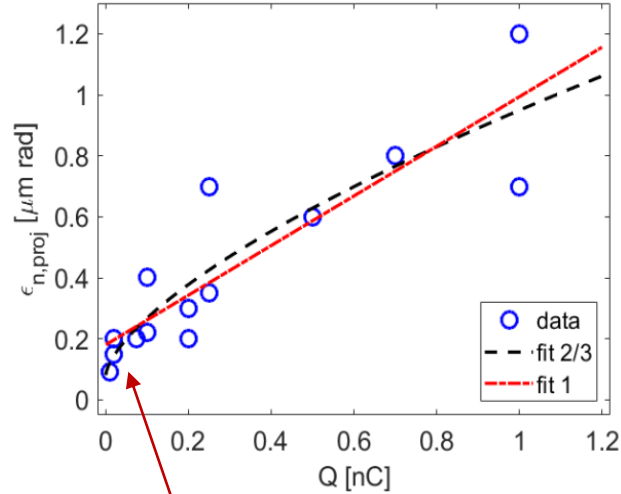
$$\varepsilon_x [\mu\text{m}] \approx b \times I [\text{kA}]$$

□ If we force  $\sigma_z = \text{const.}$  through laser shaping, we find instead:

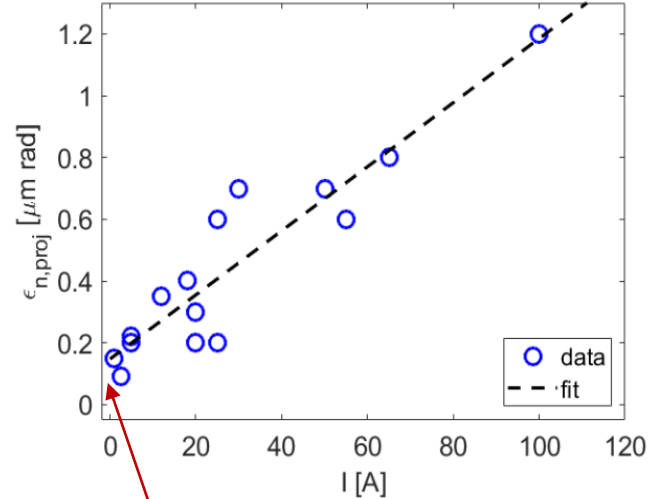
$$\varepsilon_x [\mu\text{m}] \approx c \times Q [\text{nC}]$$

# Fitting Experimental Data

Heterogeneous ensemble of experimental data from RF-PI (2009–2018, <1nC):



$\epsilon_x \sim Q^{2/3}$  is able to capture emittances of very low charges



The model has an offset for  $I \rightarrow 0$ , which takes in to account the thermal emittance

$$\epsilon_x [\mu\text{m}] \approx Q [\text{nC}]^{2/3}$$

agrees with simulations in [1] for SC-dominated beams.

Take simul. of 75pC as a ref. ("0"):

$$\sigma_{z,i} = \left(\frac{Q}{Q_0}\right)^{1/3} \sigma_{z,0}$$

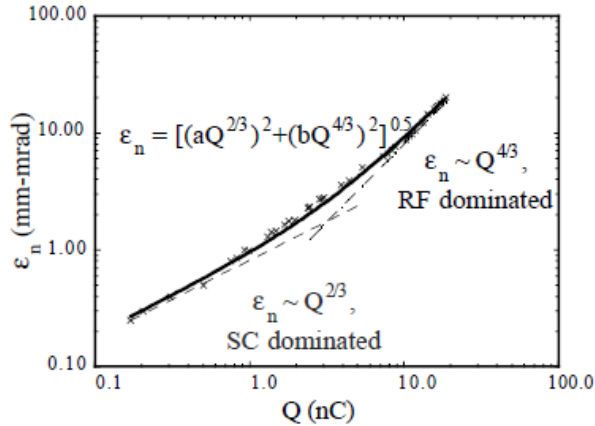
**➔**  $I_i = I_0 \left(\frac{Q}{Q_0}\right)^{2/3}$

$$CF = \frac{I}{I_i} = \frac{I}{I_0} \left(\frac{Q_0}{Q}\right)^{2/3}$$

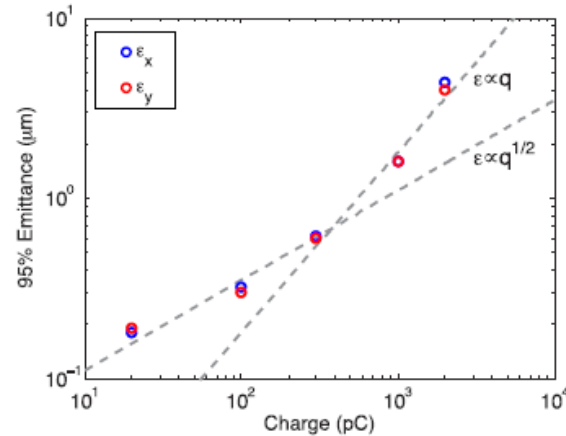
$$\sigma_{z,f} = \sigma_{z,i} / CF$$

# Examples from the literature

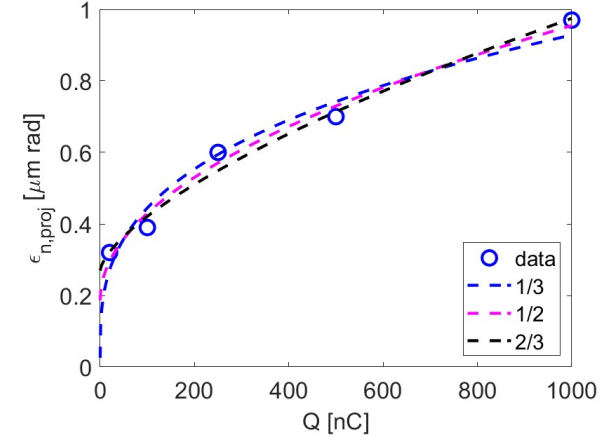
A.J. Rosenzweig and E. Colby, IEEE 1996  
*Parmela simulations of NC RF optimized Gun*



A. Bartnik et al., PRst\_AB 18, 083401 (2015)  
*Cornell DC Gun*



Data courtesy of Y. Chen for EU-XFEL at FLS'23



Most authors agree on the fact that, in SC-dominated regime, the emittance scales like  $\epsilon_{x,y} \propto Q^\nu$ , with  $\nu = \left[\frac{1}{2}, \frac{2}{3}\right]$ . For  $Q > 100\text{s pC}$ , one can observe  $\nu \rightarrow 1$ .



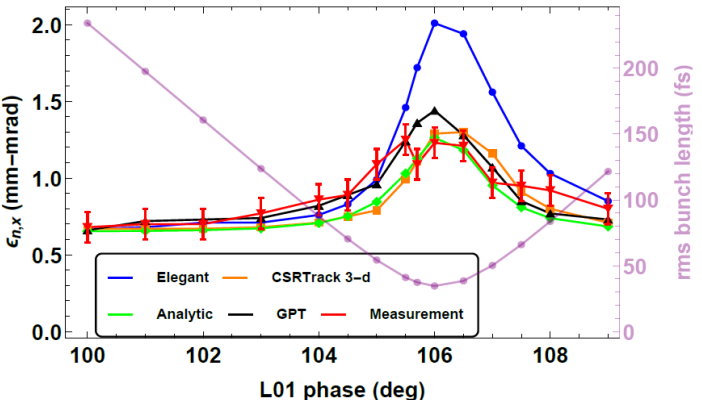
# Coherent Synchrotron Radiation

A. Brynes et al., NJP 20, 073035 (2018)  
G. Stupakov, arXiv, 1901:10745 (2019)

**The model agrees with 3-D codes. Longitudinal CSR field dominates.**

$$\Delta\epsilon_{n,L} = 7.5 \times 10^{-3} \frac{\beta}{\gamma} \left( \frac{Nr_e L_b^2}{R^3 \sigma_z^3} \right)^2$$

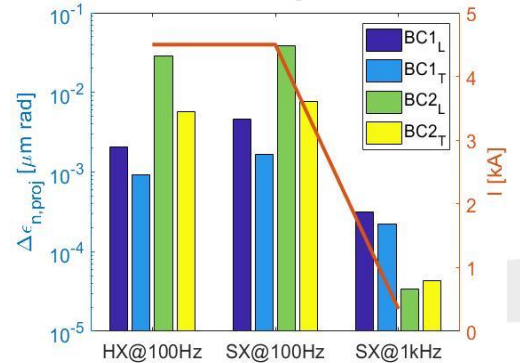
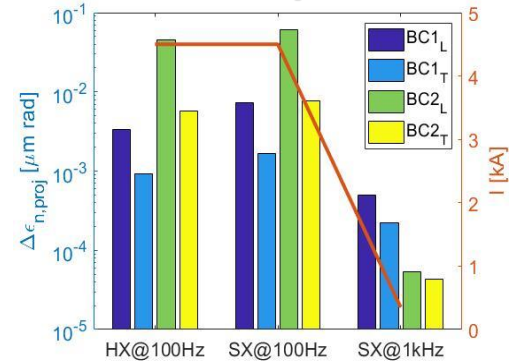
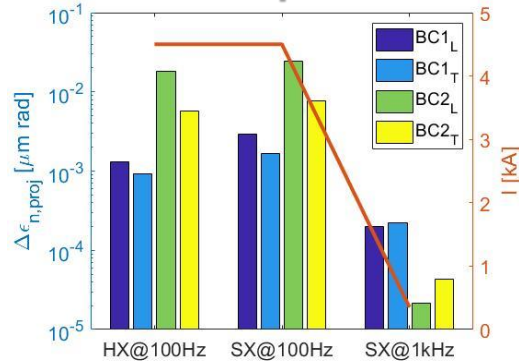
$$\Delta\epsilon_{n,T} = 2.5 \times 10^{-2} \frac{\beta}{\gamma} \left( \frac{Nr_e L_b}{R \sigma_z} \right)^2$$



**75 pC**

**150 pC**

**300 pC**







# Single Bunch Beam Break Up

$$\Delta\epsilon_{n,w_T} \propto \Delta^2 \times N_e^2 w_T^2 (2\sigma_z) \frac{L_{rf}}{\alpha G_{rf}} F(\Delta\mu) \left[ \left( \frac{\gamma_f}{\gamma_i} \right)^\alpha - 1 \right]$$

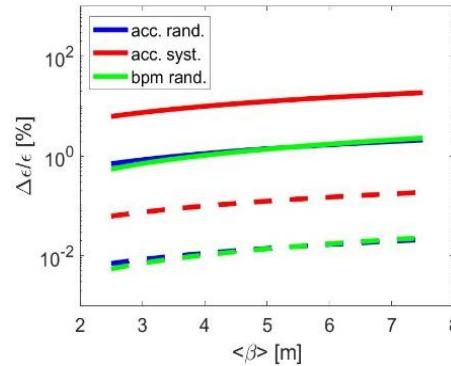
with  $\overline{\beta_u} \propto \gamma^\alpha, \alpha < 1$

1. RF random misalignment
2. RF systematic misalignment (2-by-2)
3. BPMs misalignment (FODO)

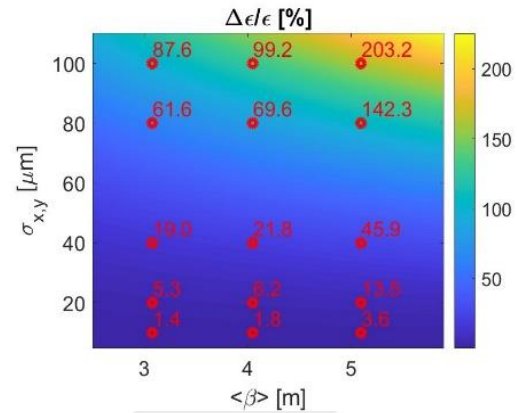
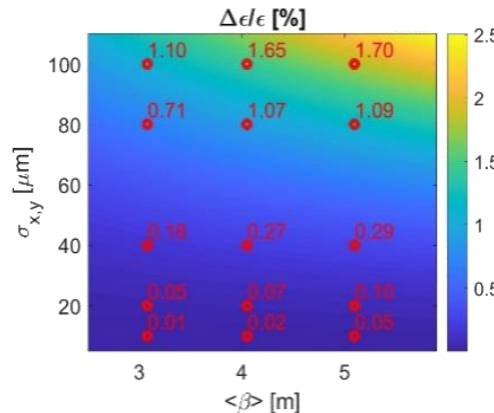
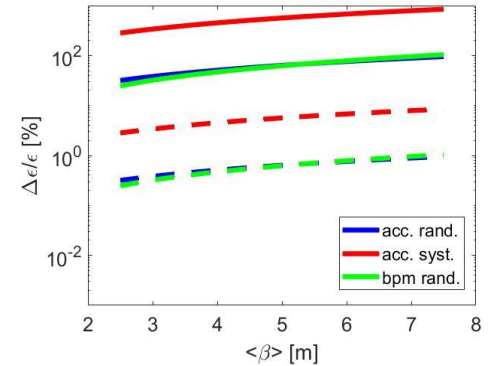
T. Raubenheimer, PRST-AB 3, 121002 (2000)

*The model agrees well with PLACET simulations. Still valid at large emittance growths.*

### 75 pC



### 300 pC





# Peak Brilliance: 3-D “slice” corrections

M. Xie, Proc. of PAC’95

G. Dattoli and M. Renieri, Laser Handbook (1986)

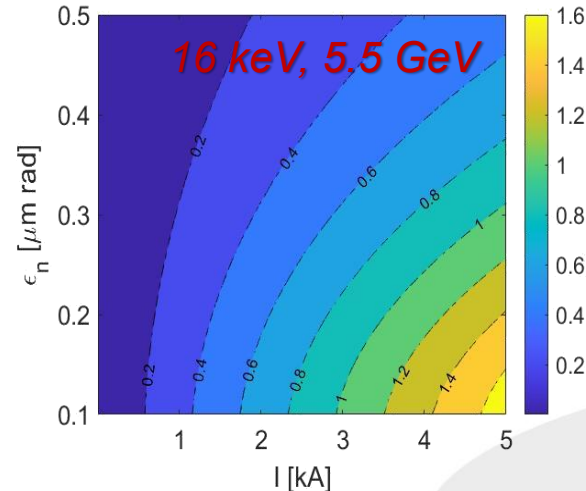
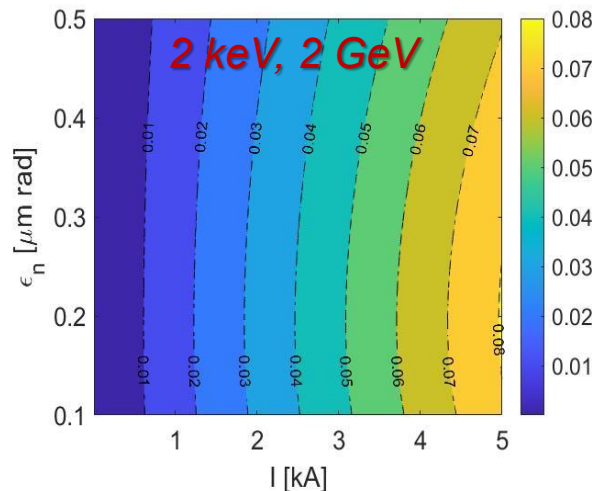
E. Saldin et al., in FEL’04 and NJP 12 (2010)

$$B_{ph} \cong 4.5 \times 10^{30} \frac{I(kA) \times E(GeV)}{\lambda(nm)} \times \delta$$

in [#ph/s/mm<sup>2</sup> /mrad<sup>2</sup> /0.1 %bw]

Sensitivity study, no functional dependence yet

- Fraction of unity
- Related to the coherent fraction of light
- **Depends on current and slice emittances through  $L_{G,3D}$**





# FEL Pulse Energy: 3-D “projected” corrections

T. Tanaka et al., NIM A 528 (2004) 172

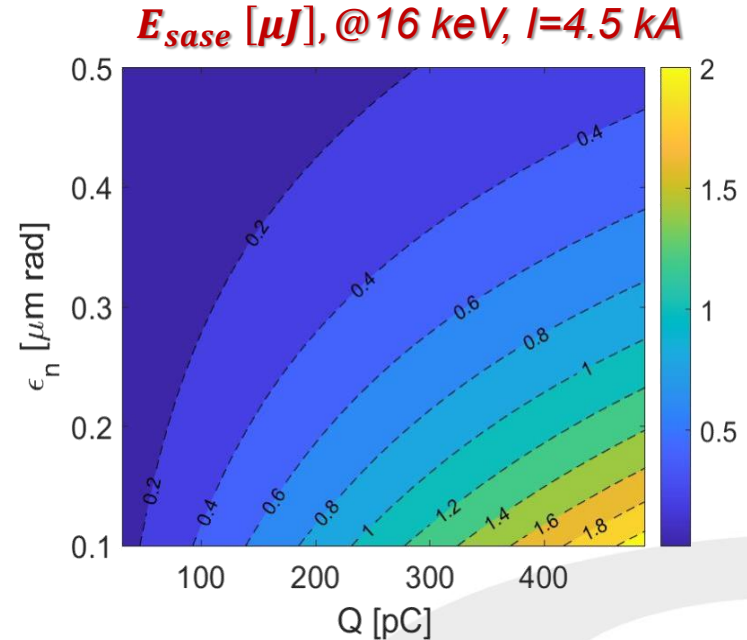
S. Di Mitri & S. Spampinati, PRST-AB 17, (2014)

$$E_{sase} \cong 0.6 \times P_{sat} \times \sqrt{2\pi\sigma_t} \cong 2.4\rho IE \frac{Q}{I} = 2.4EQ\rho$$

from spiky emission  
in t-domain from a  
flat-top bunch

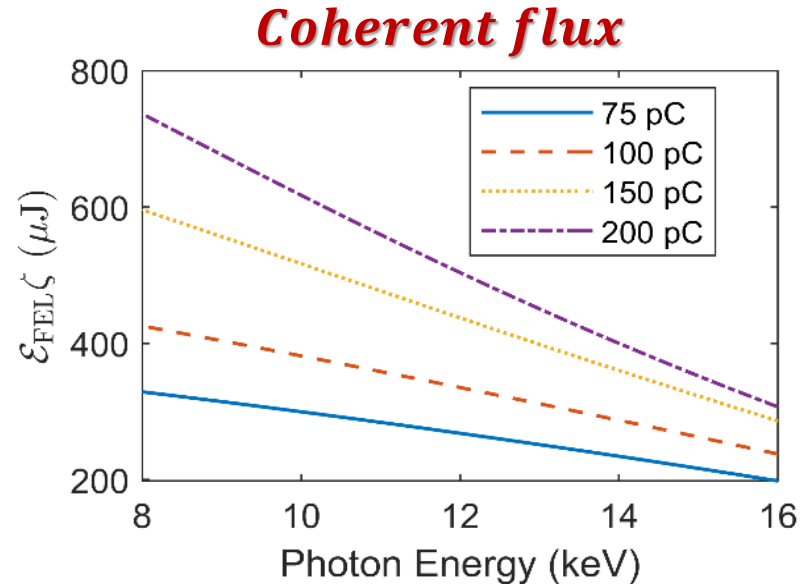
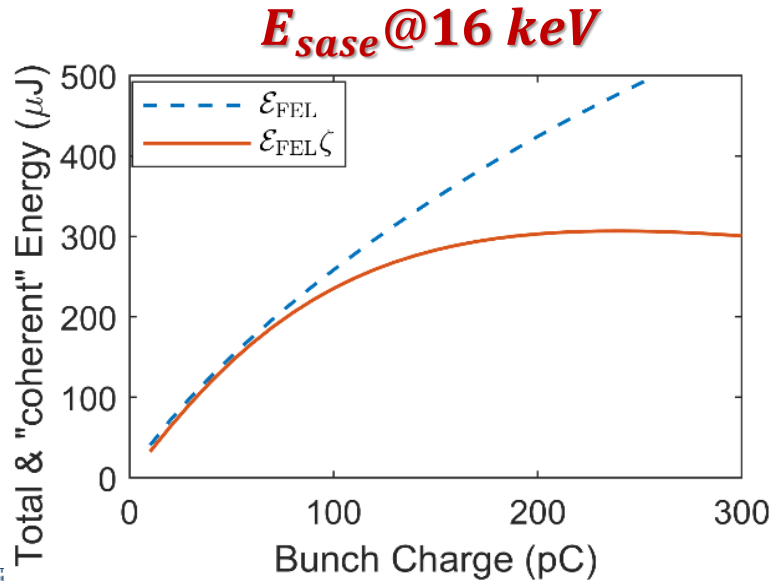
$$\rho = \frac{\rho_{3D}}{1 + \kappa \langle \theta_{coll}^2 \rangle} \approx \frac{\rho_{3D}}{1 + \left( \frac{L_{g,3D}}{2\pi\beta_u} \right) \frac{\Delta\varepsilon_{n,pr}}{\varepsilon_{n,sl}}}$$

Projected emittance growth reduces the overlap of photons and electrons during the exp amplification



# FEL Pulse Energy: Q-scaling

Finally, we apply the scaling  $\varepsilon_{n,sl} = \varepsilon_{n,sl}(Q)$ , include  $\Delta\varepsilon_{n,pr}(Q)$  from CSR and BBU, and estimate the coherent fraction of SASE flux,  $\zeta = \frac{1.1\varepsilon^{1/4}}{1+0.15\varepsilon^{9/4}}$ , with  $\varepsilon := 2\pi\varepsilon_x/\lambda$ .





## Conclusions

- ✓ The model of “invariant beam envelope” predicts space charge-emittance reasonably well for a large variety of optimized PI.
- ✓ A strategy to estimate  $E_{sase}(Q)$  is presented. 3-D slice and projected corrections to the brilliance are included in a semi-analytical, self-consistent model.
- ✓ The model is expected to highlight dominant FEL dependences from “macroscopic” e-beam parameters.



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***Thank You for Your attention***